

Contact potential difference on the atomic-scale probed by Kelvin Probe Force Microscopy: an imaging scenario

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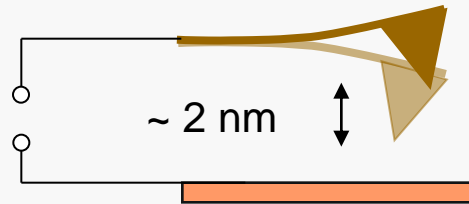
Madrid, June 17-19th 2009

Kelvin Probe Force Microscopy: concept

Goal : deriving the Contact Potential Difference of a sample w.r.t. the AFM tip

Bias voltage applied between the tip and the surface :

$$V_b = V_{DC} + V_{AC} \sin(2\pi f_{mod} t)$$



Noncontact-AFM:

Cantilever vibrates in its fundamental eigenmode (f_0)

Attractive electrostatic force:

$$F_{el} = \frac{1}{2} \frac{\partial C}{\partial z} V^2$$

with

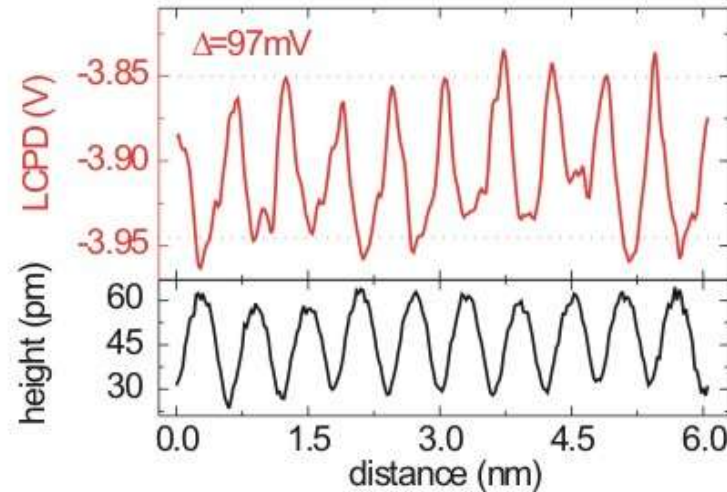
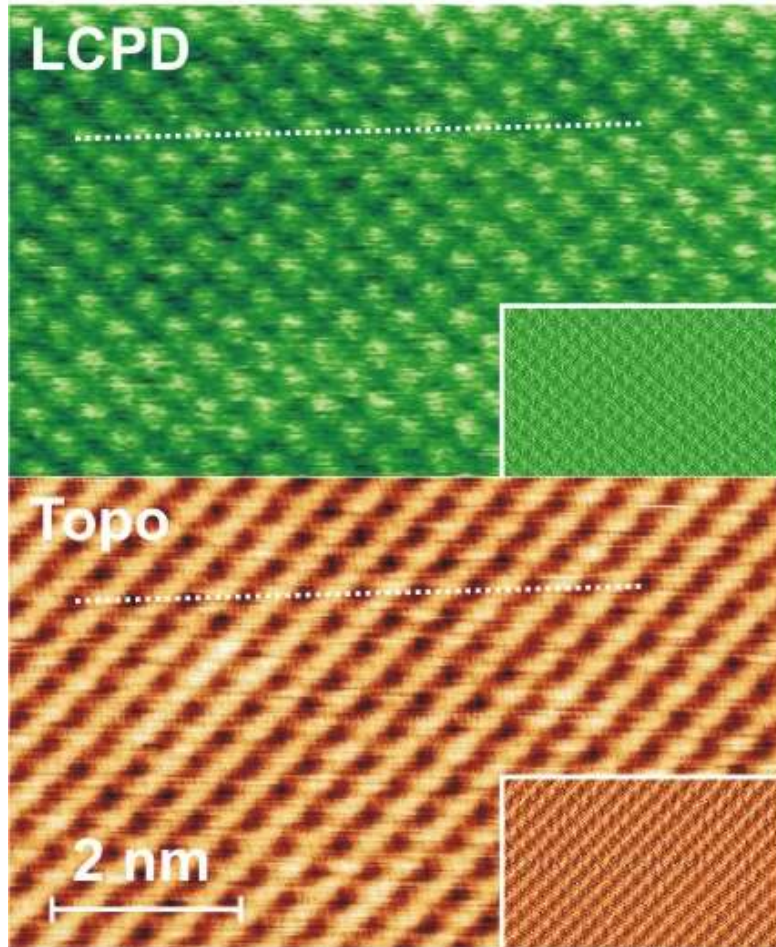
$$V = V_b - V_{CPD} = V_b - \Delta\Phi / e$$

$\Delta\Phi$ = Work functions difference between tip and sample (= CPD)

- ⇒ Induced vibrations of the cantilever at f_{mod} and $2f_{mod}$
- ⇒ Detection and cancellation of the f_{mod} component (Δf in FM-KPFM, $A^{(1)}$ in AM-KPFM) by applying a proper DC bias voltage which matches the CPD
- ⇒ **V^2 DEPENDENCE OF THE FORCE REQUIRED, OTHERWISE NO KPFM CONTRAST**

Atomic-scale KPFM contrast on ionic surfaces

AM-KPFM (UHV) on bulk KBr(001)¹

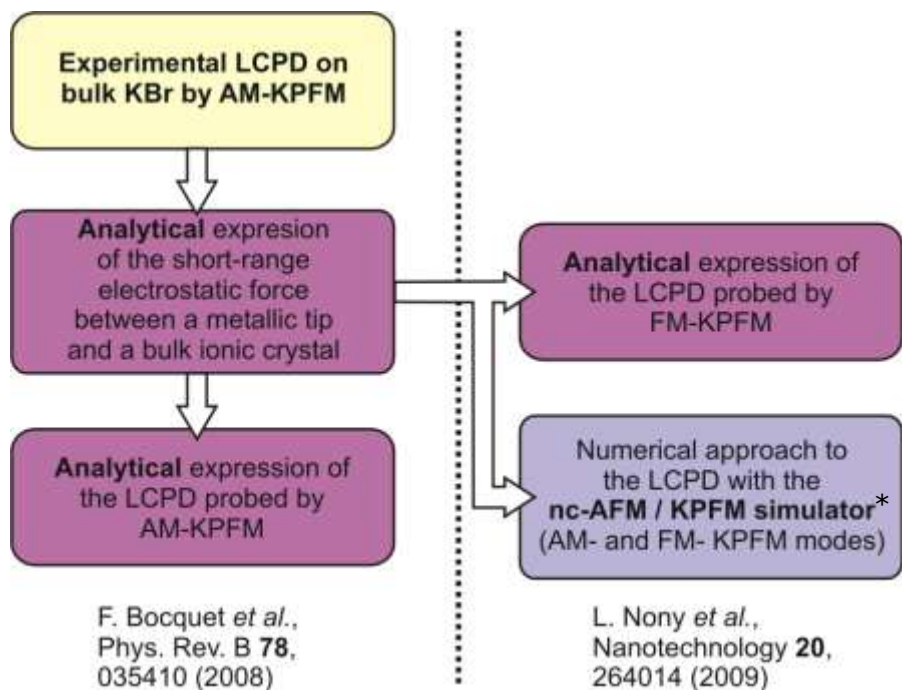


- ❑ Atomic-scale contrast of the CPD (Local CPD, LCPD):
 - ✓ Madelung surface potential of the crystal?
 - ✓ Compliant with the magnitude of the measured signal?
 - ✓ Instrumentation (controllers)?
- ❑ Atomic-scale KPFM contrast as well reported on SC samples (earlier results), but neither understood in details

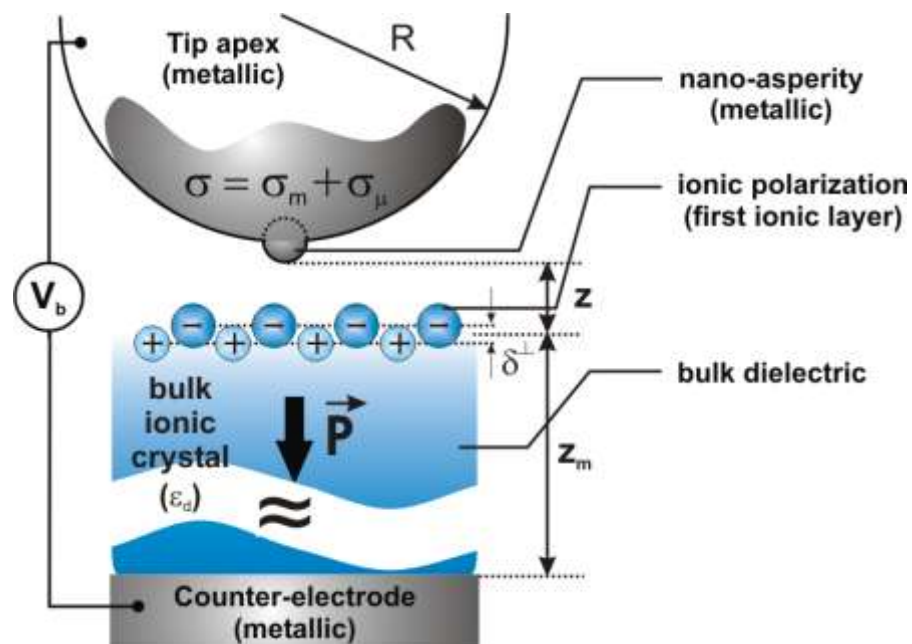
¹F. Bocquet *et al.*, Phys. Rev. B **78**, 035410 (2008)

- **Short-Range and bias-dependent Electrostatic forces (SRE forces) play a key role in the occurrence of the atomic-scale contrast in KPFM images.**
- **What are the underlying processes?**
- **How do they influence the experimental measurement of the LCPD in terms of magnitude and lateral resolution?**

Our approach to the LCPD on KBr so far



Geometry of the problem:

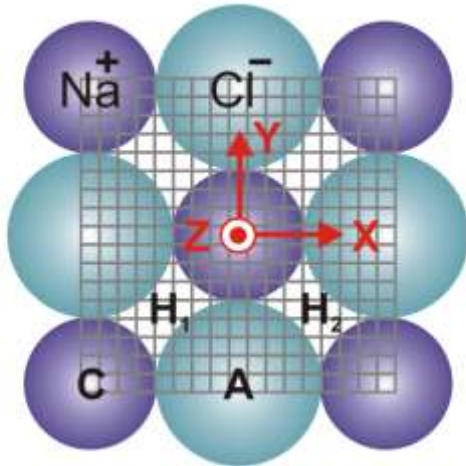
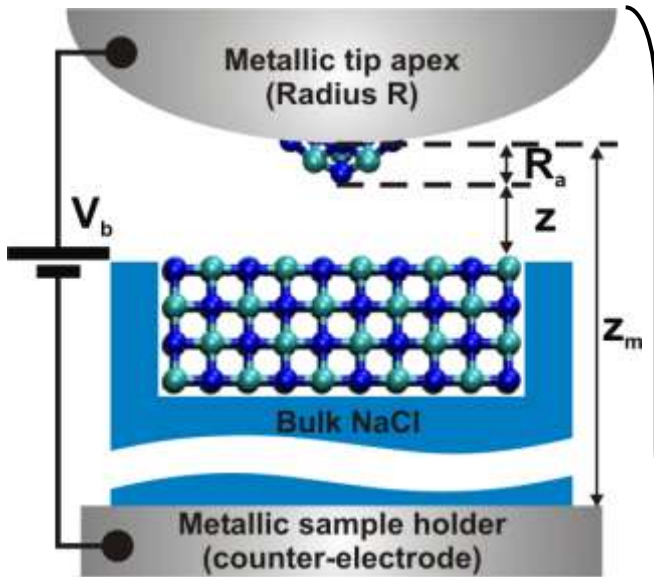


Major findings:

- ❑ SRE forces, i.e. bias dependent forces with a short-range nature, do occur between the AFM tip and the surface when they are
 - ❑ Not or
 - ❑ The q
 - ❑ surface
 - ❑ The L
 - BUT not
- Combination between analytical SRE force and the simulator: useful to understand the physics of the process. However, owing to the unrealistic geometry of the tip, large LCPD values are predicted (+/- 5 V between ionic and cationic sites) → not fully consistent with the experimental results (ranging from 0.1V up to +/- 2 V maximum)
- BETTER DESCRIPTION OF THE TIP REQUIRED!**

*L.Nony *et al.* Phys. Rev. B **74**, 235439 (2006)

Improving the model: fully numerical approach



L.Nony *et al.*, accepted in Phys. Rev. Lett.

□ Atomistic simulations (SciFi):

- ✓ ions treated atomistically in a core-shell model
- ✓ **ionic relaxation and electronic polarization and ionic polarization are taken into account**
- ✓ force convergence criterion : 1 meV / Å / ion

□ Tip:

- ✓ metallic tip apex : 5nm radius
- ✓ tip **cluster**: 4x4x4 described atomistically; oriented along the [111] direction w.r.t. the surface, Na-terminated
- ✓ protrudes with an height $R_a = 3\text{\AA}$ from the metallic apex: **carries a charge of +1 compensated by applying -0.91V to the tip**

□ Sample:

- ✓ NaCl **slab**: 10x10x4 described atomistically
- ✓ bulk NaCl: 5mm thick, described as a continuous medium

□ Mesh of the unit cell:

- ✓ $X \times Y \times Z \times V = 17 \times 17 \times (8\text{nm}-0.3\text{nm})/0.005\text{nm} \times (2.3\text{V} - -3.4\text{V})/0.01 \text{ V} \sim 255 \text{ Msamples}$ to handle
- ✓ **four particular sites investigated: cationic (C), anionic (A) and inequivalent hollow (H_1, H_2) sites**

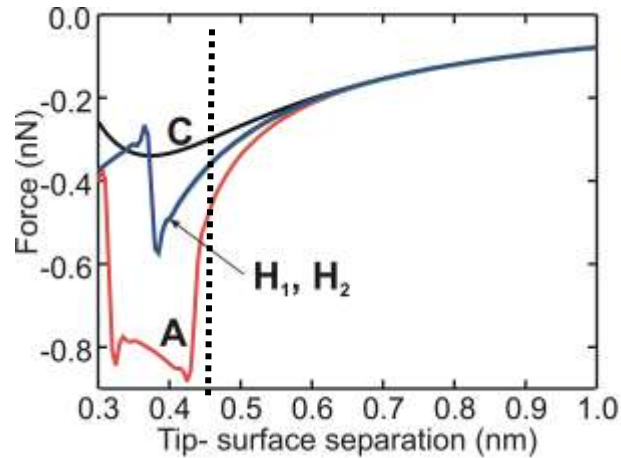
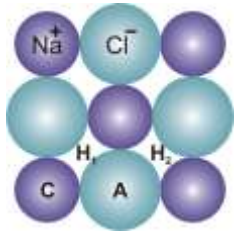
(X, Y, Z, Bias)
4D force field
+ VdW
+ long-range
capacitive force ($\sim V_b^2$)

nc-AFM / KPFM
simulator
(FM-KPFM mode)

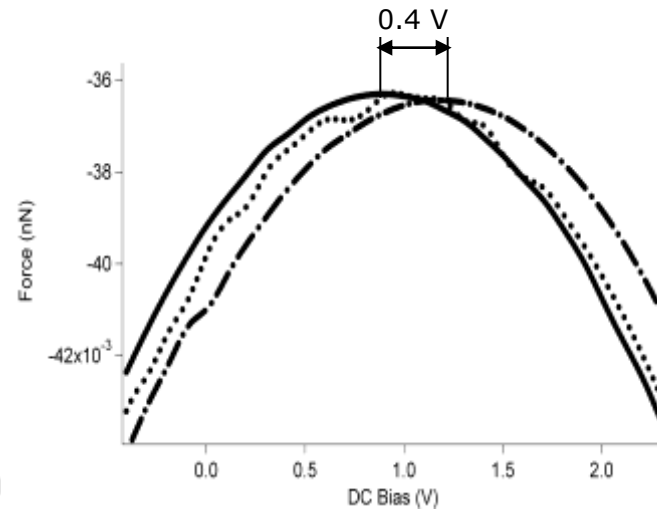
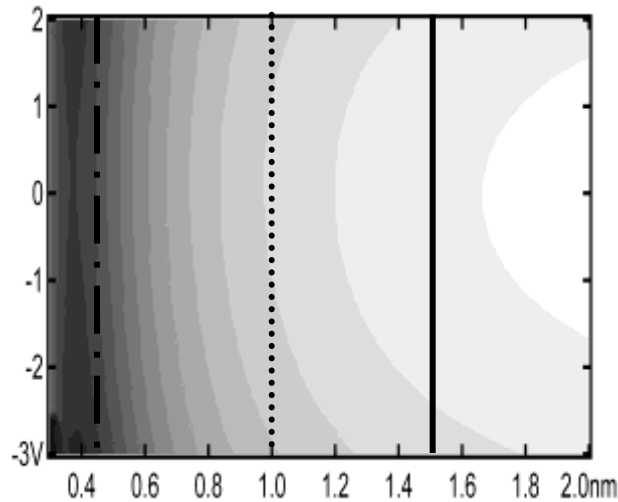
Topographical
& CPD images

Spectroscopic
curves

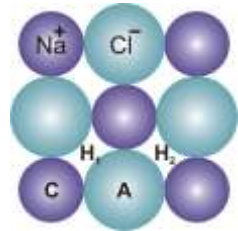
Force field: evidence for SRE forces



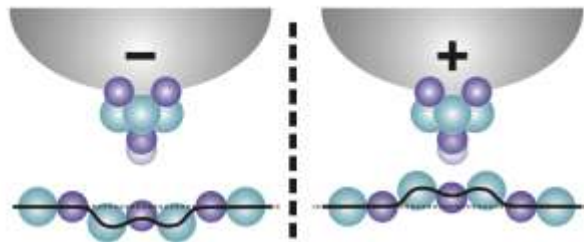
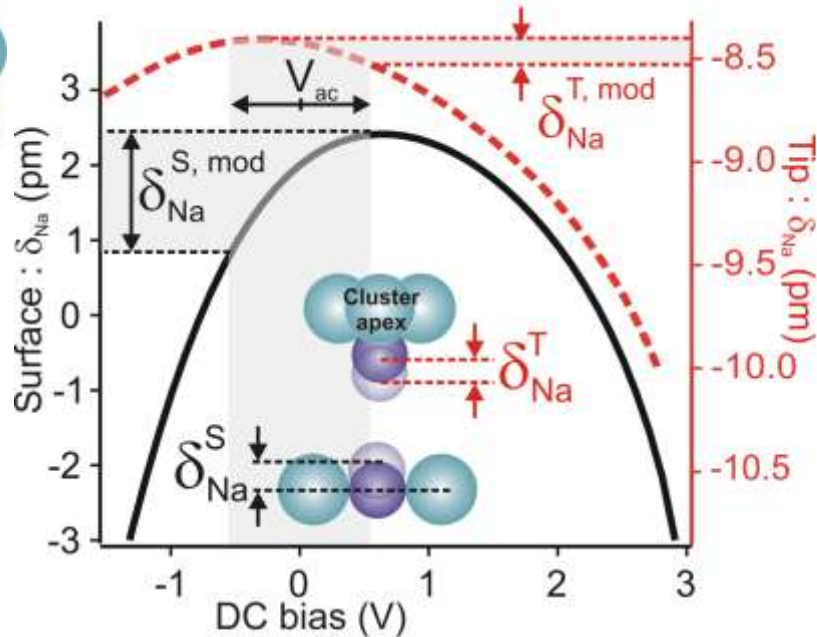
Evidence for SRE forces (top of a cation)



Ionic displacements at the tip-surface interface

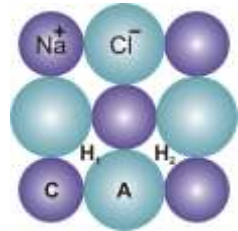


Cores displacements of the closest ions at the interface:



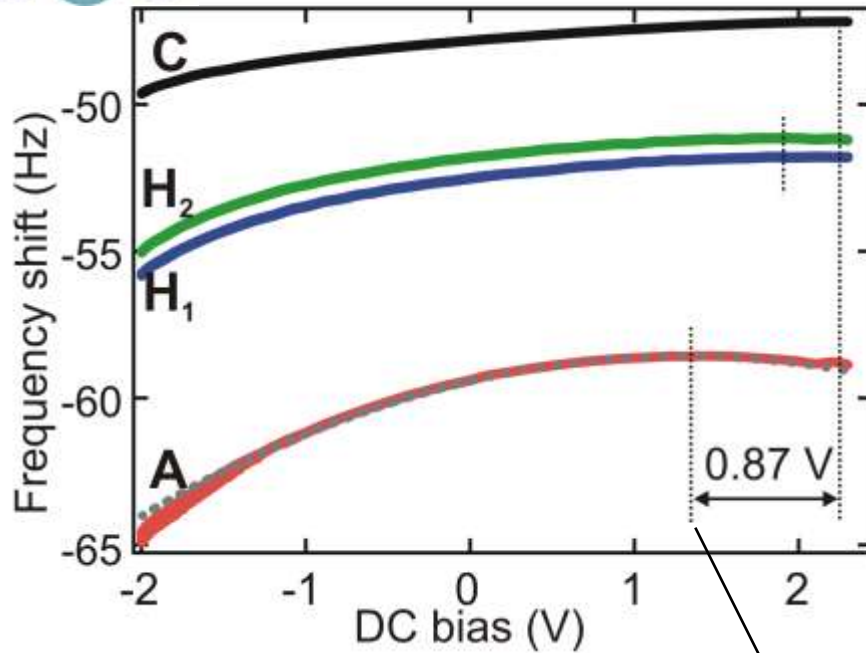
- AC Modulation of the applied bias triggers ionic displacements at the tip-surface interface
- Self-consistent influence between chemical and SRE forces
- Deviation from the usual capacitive (parabolic-like) force

Connection with the LCPD by means of the nc-AFM/KPFM (FM-KPFM) simulator



Simulated spectroscopic curves (Δf vs. V_{DC}):

$z = 0.45 \text{ nm}$



-LCPD @ $z=0.45\text{nm}$ site A

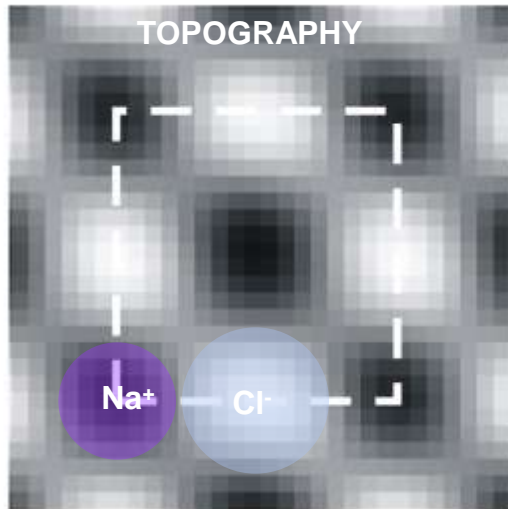
- No parabolic dependence (consistent with the force)
- LCPD is site-dependent (atomic-scale contrast possible)
- Resonance effect: experimentally observed and theoretically predicted

Numerical topographical and KPFM images

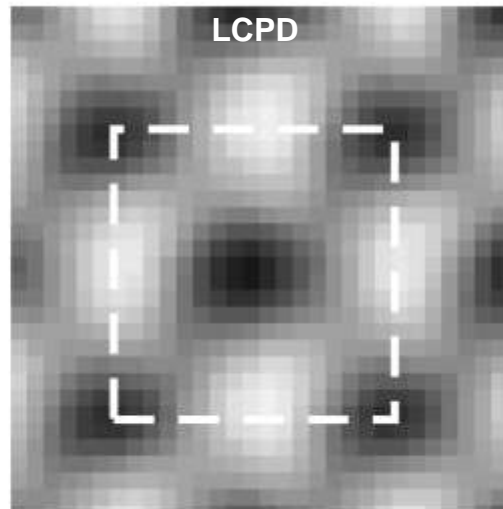
Simulated scans:

- nc-AFM (main) parameters: $A_0 = 8$ nm p-p ; $Q = 35000$; $f_0 = 150$ kHz
- FM-KPFM parameters: $V_{AC} = 0.5$ V ; $f_{mod} = 1$ kHz
- scan speed : 1.5 s/line

Distance controller engaged: $\Delta f = -47.22$ Hz ($z_{start} \sim 0.45$ nm)

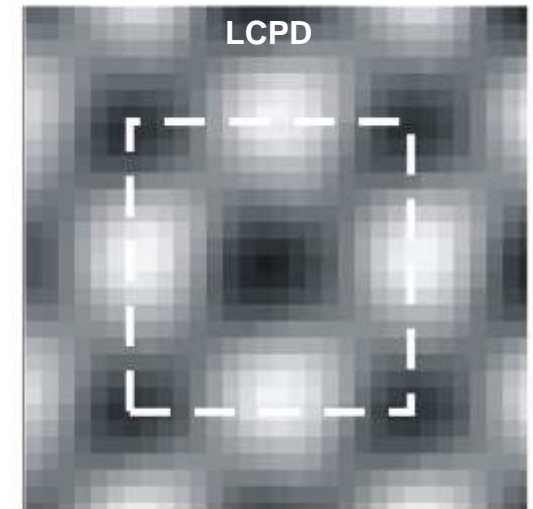


Contrast range: 38 pm



Contrast range: 0.56 V

Constant height: $z = 0.45$ nm



Contrast range: **0.86 V**
→ consistent with
the spectro. curves

- Simultaneous atomic-scale topographical and KPFM contrast
- Magnitudes of the contrasts **now consistent with the experimental observations** (30pm, 0.1V)
- The z-regulation prevents from measuring the “correct” LCPD value (spectro. curves)

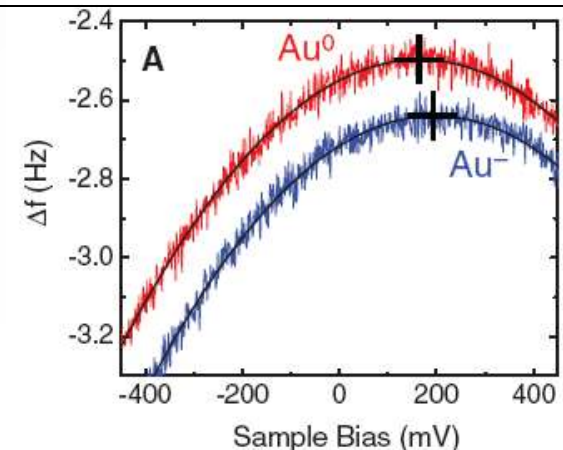
Take-home message

- Fully numerical approach to the atomic-scale KPFM contrast: combination of atomistic simulations with the nc-AFM / KPFM simulator
- SRE are site-dependent, which makes the achievement of the atomic-scale KPFM contrast possible
- The atomic-scale KPFM contrast relies on the ionic displacements at the tip-surface interface that are triggered by the AC modulation of the applied bias (in the range of ~ 1 to 2 pm only!!!)
- The LCPD contrast has the lateral periodicity of the Madelung surface potential but does not match it owing to self-consistent influence of chemical and SRE forces

Measuring the Charge State of an Adatom with Noncontact Atomic Force Microscopy

Leo Gross,^{1*} Fabian Mohn,¹ Peter Liljeroth,^{1,2} Jascha Repp,^{1,3} Franz J. Giessibl,³ Gerhard Meyer¹

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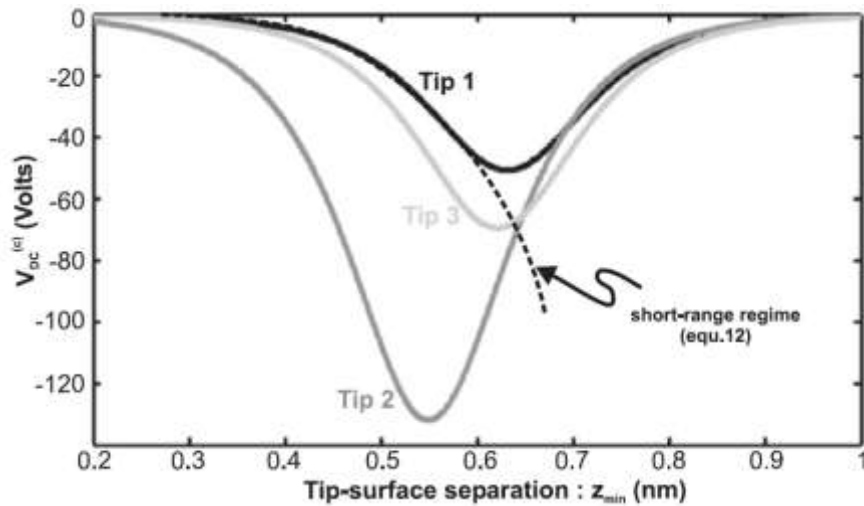
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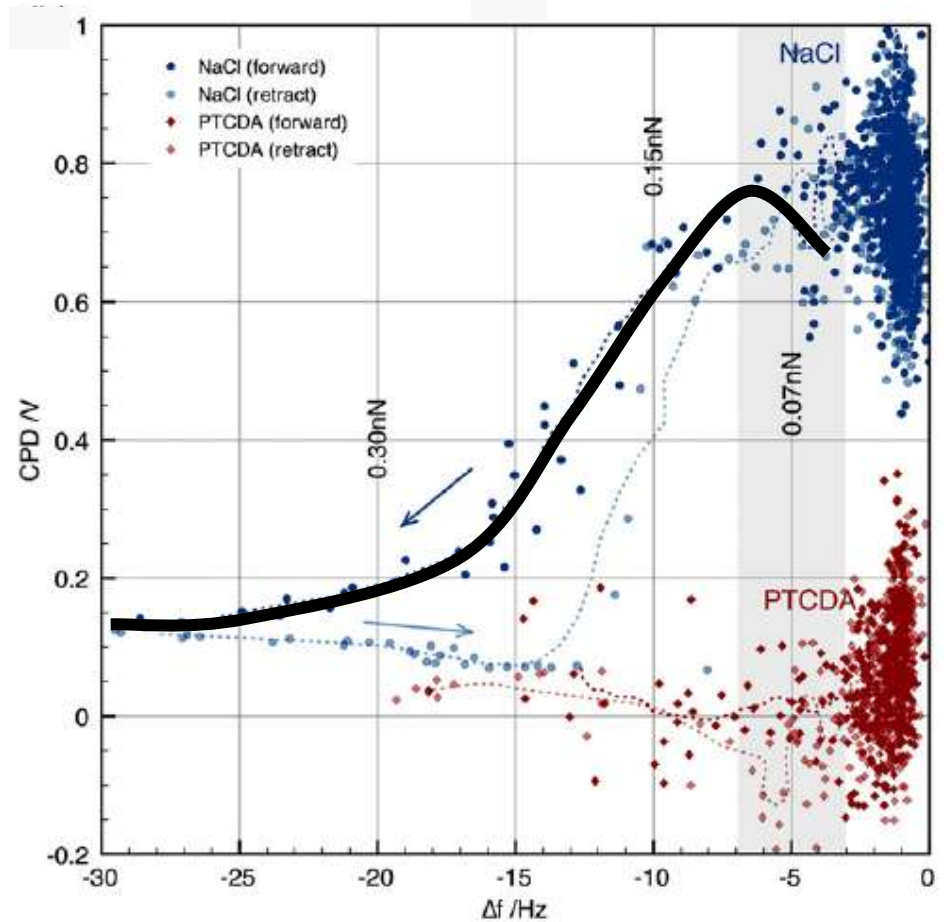
T. Glatzel

S. Kawai

Resonance effect of the LCPD



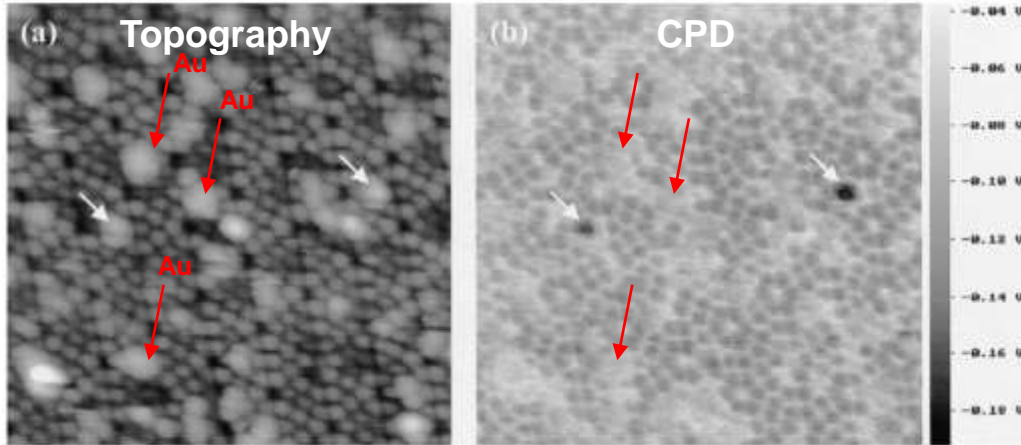
L. Nony et al., Nanotechnology **20**, 264014 (2009)



S. Burke et al., Nanotechnology **20**, 264012 (2009)

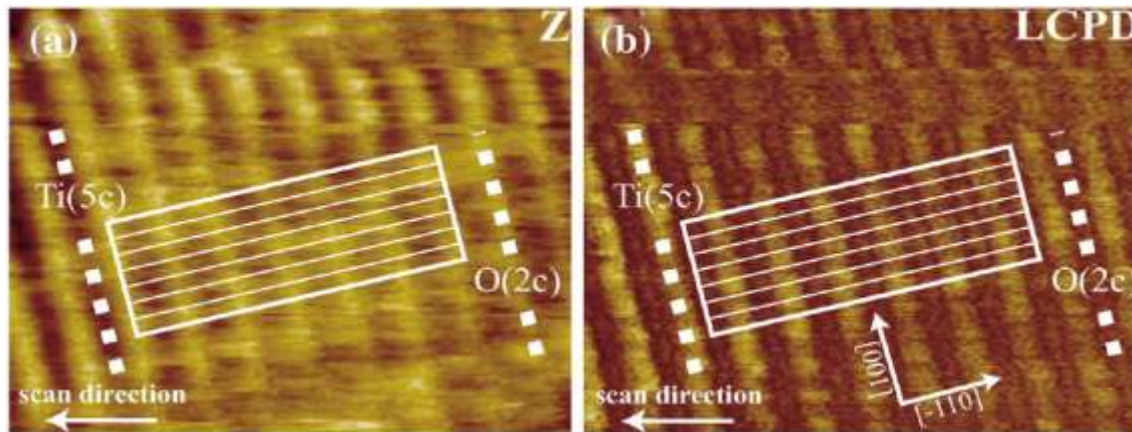
High-resolution KPFM: what is being measured?

- S. Kitamura *et al.*, Appl. Surf. Sci. **157**, 222 (2000) : Au/Si(111) 7x7



"The potential difference of atomic structures...**does not seem to reflect the work function** as we initially expected. It is therefore considered that the atomic potential difference reflects the local electron density on the surface."

- G.H. Enevoldsen *et al.*, Phys. Rev. Lett. **100**, 236104 (2008) : TiO₂(110)



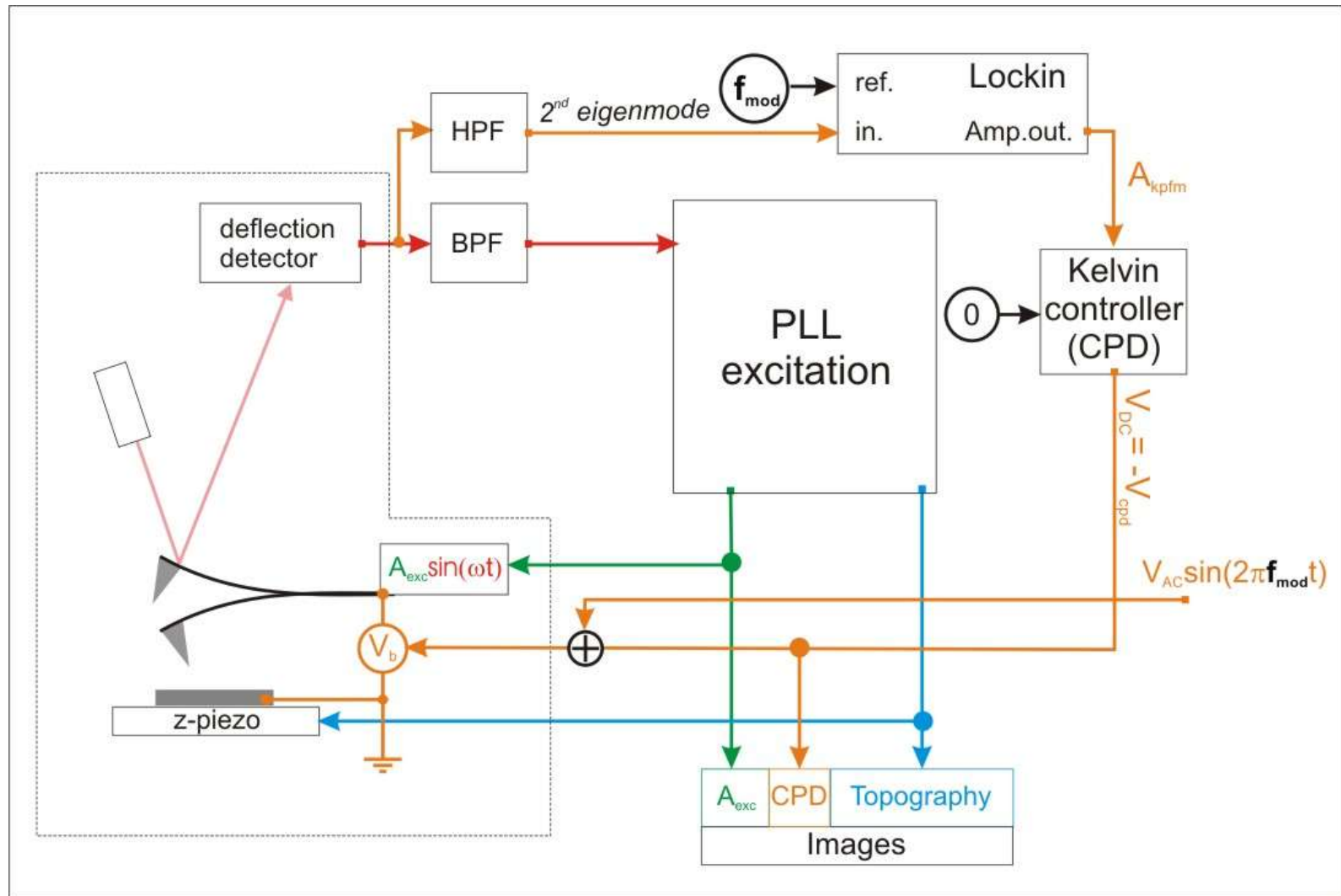
"... the absolute values of the LCPD are determined to a very large degree by the tip sharpness and the conductive nature of the tip (metallic, ionic), and a quantitative comparison between atomic-scale LCPD measurements and surface potentials derived by theory should take this finding into account."

Features of the nc-AFM simulator¹

- ❑ Numerical implementation of a real nc-AFM electronic setup² (mix of analog/digital devices)
- ❑ The Phase-Locked Loop (PLL) detects Δf AND provides the phase-shifted excitation to the cantilever
- ❑ Initially developed to address the issue of *apparent dissipation*
- ❑ Main results :
 - No unwanted phase-shift if the PLL center frequency is continuously updated : *i.e.* no apparent dissipation
 - No apparent dissip. while : $2 \times \tau_{\text{PLL}} \sim 0.7 \text{ ms} < \tau_{\text{APIC}} \sim 2 \text{ ms} < \tau_{\text{DPIC}}$
 - Apparent dissipation can conditionally be generated, but can't exceed 15% of the intrinsic dissipation of the free cantilever : likely not experimentally relevant in UHV at room temp.

¹L. Nony *et al.*, Phys. Rev. B **74**, 235439 (2006), ²Ch. Loppacher *et al.*, Appl. Phys. A **66**, 215 (1998)

Numerical implementation of the **AM-KPFM** mode within the nc-AFM simulator



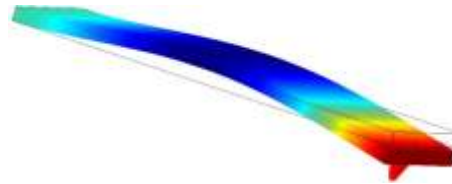
Numerical description of the coupling between the eigenmodes in **AM-KPFM**

First eigenmode :

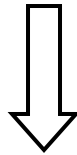


$f_0 \sim 150$ kHz
 $k_0 \sim 30$ N/m
 $Q_0 \sim 30000$ ($\Delta f_{\text{HWHM}} = 2.5$ Hz)
 Actuated mechanically

Second eigenmode :



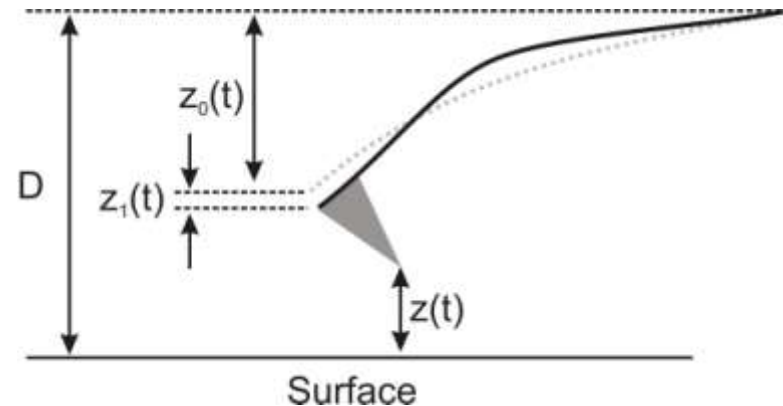
$f_1 = 6.3f_0 \sim 945$ kHz
 $k_1 \sim 36k_0 \sim 1100$ N/m
 $Q_1 \sim 8000$ ($\Delta f_{\text{HWHM}} = 60$ Hz)
 Actuated electrostatically



Coupling occurs due to the **interaction force**

Instantaneous tip-surface distance :

$$z(t) = D - z_0(t) - z_1(t)$$

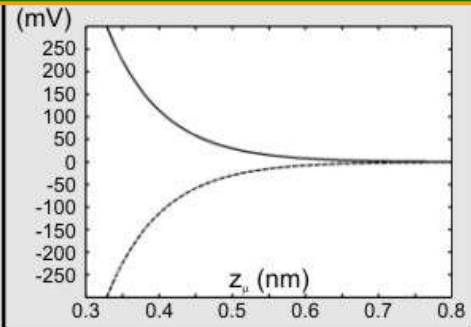
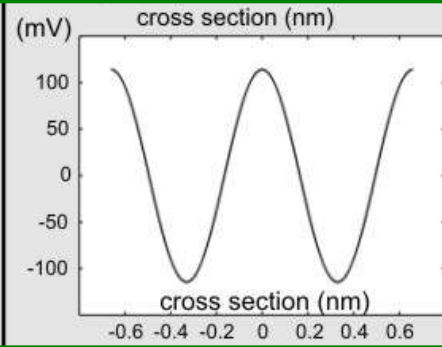
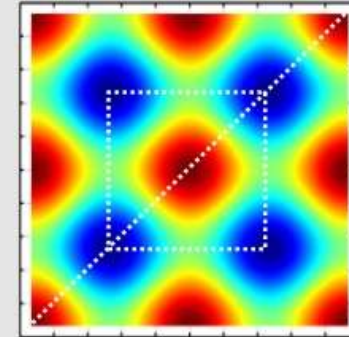


How is the Madelung potential derived? Why is one single topmost ionic layer enough?

Madelung surface potential¹:

$$V_s(x, y, z_\mu) = - \frac{q}{\pi \epsilon_0 a'} \cosh[\tilde{\delta}^\perp(V_b)] \tilde{\chi}(x, y) e^{-(2\pi/a')z_\mu}$$

Madelung surface potential' on KBr @ $z_\mu = 4 \text{ \AA}$



¹ R.Watson *et al.*, Phys. Rev. B **24**, 1791 (1981)